#### **Ehrenfest dynamics**

## $\epsilon(\mathbf{R}) ightarrow \langle \tilde{\Phi}(\mathbf{R}^{lpha}(t),t) | \hat{\mathcal{H}}_{el}(\mathbf{R}^{lpha}(t)) | \tilde{\Phi}(\mathbf{R}^{lpha}(t),t) angle_{\mathbf{r}}$



Credit: J. Michl and V. Bonacic-Koutecký, "Electronic Aspects of Organic Photochemistry". John Wiley and Sons, Inc., New York, 1990.

#### Ehrenfest dynamics

In Ehrenfest dynamics, the potential energy  $\epsilon(\mathbf{R})$  used to evolve a trajectory is the average electronic energy as computed from the time-dependent electronic wavefunction:

Ehrenfest dynamics

$$i\hbar\dot{c}^{\alpha}_{I}(t) = \sum_{J}c^{\alpha}_{J}(t)\left(E^{eI}_{J}(\mathbf{R}^{\alpha}(t))\delta_{IJ} - i\hbar\dot{\mathbf{R}}^{\alpha}\cdot\mathbf{d}_{IJ}(\mathbf{R}^{\alpha}(t))\right)$$

$$\begin{aligned} \mathsf{F}^{\alpha} &= -\nabla_{\mathsf{R}} \langle \tilde{\Phi}(\mathsf{R}^{\alpha}(t),t) | \hat{\mathcal{H}}_{el}(\mathsf{R}^{\alpha}(t)) | \tilde{\Phi}(\mathsf{R}^{\alpha}(t),t) \rangle_{\mathsf{r}} \\ &= -\nabla_{\mathsf{R}} \left( \sum_{IJ} c_{I}^{\alpha}(t)^{*} c_{J}^{\alpha}(t) \langle \Phi_{I}(\mathsf{R}^{\alpha}(t)) | \hat{\mathcal{H}}_{el}(\mathsf{R}^{\alpha}(t)) | \Phi_{J}(\mathsf{R}^{\alpha}(t)) \rangle_{\mathsf{r}} \right) \end{aligned}$$

One can perform Ehrenfest dynamics by using the time-dependent Kohn-Sham equation instead of the time-dependent electronic Schrödinger equation!!!

#### Nuclear dynamics with a mean potential



$$i\hbar\dot{c}^{\alpha}_{I}(t) = \sum_{J} c^{\alpha}_{J}(t) \left( E^{el}_{J}(\mathbf{R}^{\alpha}(t))\delta_{IJ} - i\hbar\dot{\mathbf{R}}^{\alpha} \cdot \mathbf{d}_{IJ}(\mathbf{R}^{\alpha}(t)) \right)$$

$$\mathbf{F}^{\alpha} = -\nabla_{\mathbf{R}} \left( \sum_{IJ} c_{I}^{\alpha}(t)^{*} c_{J}^{\alpha}(t) \langle \Phi_{I}(\mathbf{R}^{\alpha}(t)) | \hat{\mathcal{H}}_{eI}(\mathbf{R}^{\alpha}(t)) | \Phi_{J}(\mathbf{R}^{\alpha}(t)) \rangle_{\mathbf{r}} \right)$$

Nonadiabatic (Molecular) Dynamics

#### Nuclear dynamics with a mean potential



- Approximate account for nonadiabatic effects mean-field dynamics.
- Produce physical nonadiabatic dynamics only if the different electronic states important for the dynamics have similar potential energy surfaces, or for short-time dynamics.

# An example of Ehrenfest Dynamics dynamics using TDDFT

#### Early mechanisms of hadrontherapy



Credit: J. Michl and V. Bonacic-Koutecký, "Electronic Aspects of Organic Photochemistry". John Wiley and Sons, Inc., New York, 1990.

#### Hadrontherapy

"Process of irradiating a malignant tumor with ionizing radiation with the aim of damaging the DNA of the cancerous cells."

- When high-energy photons are used: X-ray therapy.
- When hadronic matter is used (H or C): Hadrontherapy.



#### Hadrontherapy

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#### Hadrontherapy - The idea

The main advantage of heavily-charged ion beams is their favorable depth-dose profile, also known as Bragg curve <sup>3</sup>.



<sup>3</sup>Rev. Mod. Phys. **82**, 383 (2010).

#### Setup

#### Ehrenfest dynamics with TDDFT

- Guanine-Cytosine base pair + 107 water molecules.
- Full system treated at the Ehrenfest/TDDFT level of theory (PBE functional).
- Projectile: C<sup>4+</sup>.
- Speed of the ion in the Bragg peak:  $\sim 0.79\%$  of c  $\Rightarrow$  relativistic effects are negligible!
- Projectile added on a corner of the box, with a velocity computed to cross the bubble with no direct collision.
- Time step = 0.24 as. Total duration of the dynamics: 1 fs.
- And the final setup is...

Unpublished results.

#### Setup



Application of Ehrenfest dynamics to hadrontherapy

## Electron extraction from C<sup>4+</sup> - [ $\rho(\mathbf{r}, t)$ ]





Application of Ehrenfest dynamics to hadrontherapy

Electron extraction from C<sup>4+</sup> -  $\left[\rho(\mathbf{r}, t) - \rho_0^{(\text{no C}^{4+})}(\mathbf{r})\right]$ 



## **Trajectory Surface Hopping dynamics**

## $\epsilon(\mathbf{R}) ightarrow E^{el}_*(\mathbf{R})$



Credit: J. Michl and V. Bonacic-Koutecký, "Electronic Aspects of Organic Photochemistry". John Wiley and Sons, Inc., New York, 1990.

#### Trajectory Surface Hopping dynamics

Trajectory Surface Hopping (TSH) invokes that the PES,  $E_*^{el}(\mathbf{R})$ , from which the nuclear forces are computed can be chosen based on a set of rules after each nuclear integration time step:

Trajectory Surface Hopping dynamics

$$i\hbar \dot{c}^{\alpha}_{I}(t) = \sum_{J} c^{\alpha}_{J}(t) \left( E^{eI}_{J}(\mathbf{R}^{\alpha}(t))\delta_{IJ} - i\hbar \dot{\mathbf{R}}^{\alpha}(t) \cdot \mathbf{d}_{IJ}(\mathbf{R}^{\alpha}(t)) \right)$$

$$\mathbf{F}^{\alpha} = -\nabla_{\mathbf{R}} E_*^{el}(\mathbf{R})|_{\mathbf{R} = \mathbf{R}^{\alpha}(t)}$$

In TSH, a trajectory  $\alpha$  evolves adiabatically on a given potential energy surface but can *hop* after each nuclear time step if nonadiabatic effects are detected.

















#### What is the prescription for a hop to occur?

Tully devised in 1990 the *fewest switches* algorithm, which provides a possible prescription for a trajectory to hop from one state to the other after each time step:

$$\mathcal{P}_{K \to I}^{\alpha} = \max\left[0, \frac{-2 dt}{\left|c_{K}^{\alpha}(t)\right|^{2}} \Re\left[c_{K}^{\alpha*}(t)c_{I}^{\alpha}(t)\right] \dot{\mathbf{R}}^{\alpha}(t) \cdot \mathbf{d}_{IK}\left(\mathbf{R}^{\alpha}(t)\right)\right].$$

A hop occurs from electronic K to another electronic state I if

$$\sum_{J}^{I-1} \mathcal{P}^{\alpha}_{K \to J} \leq \zeta \leq \sum_{J}^{I} \mathcal{P}^{\alpha}_{K \to J},$$

where  $\zeta$  is a random number generated uniformly in the interval [0 : 1]. If a hop occurs, the nuclear forces will now be obtained from  $E_*^{el}(\mathbf{R}) = E_l^{el}(\mathbf{R})$ .

An important point of the TSH algorithm is that it requires a large number of *independent* trajectories, or independent TSH runs, to converge the hopping algorithm, as well as the sampling of initial conditions.

#### Trajectory Surface Hopping dynamics – central equations

Trajectory Surface Hopping dynamics

$$i\hbar\dot{c}^{\alpha}_{I}(t) = \sum_{J} c^{\alpha}_{J}(t) \left( E^{el}_{J}(\mathbf{R}^{\alpha}(t))\delta_{IJ} - i\hbar\dot{\mathbf{R}}^{\alpha}(t)\cdot\mathbf{d}_{IJ}(\mathbf{R}^{\alpha}(t)) \right)$$

$$\mathbf{F}^{lpha} = -
abla_{\mathbf{R}} E^{el}_{*}(\mathbf{R})|_{\mathbf{R}=\mathbf{R}^{lpha}(t)}$$

$$\sum_{J}^{I-1} \mathcal{P}_{K\to J}^{\alpha} \leq \zeta \leq \sum_{J}^{I} \mathcal{P}_{K\to J}^{\alpha},$$

Multiple independent trajectories  $\alpha$  have to be propagated.

#### Trajectory Surface Hopping dynamics – some warnings



• Rescaling of the nuclei velocities after a hop is still a matter of debate.

- Depending on the system studied, a lot of trajectories could be needed to obtain a complete statistical description of the non-radiative relaxation.
- TSH cannot adequately describe the decoherence of nuclear wavepackets...

For more details (and warnings) about Tully's surface hopping, see G. Granucci and M. Persico, *J. Chem. Phys.*, **126**, 134114 (2007).

## An example of Trajectory Surface Hopping dynamics using LR-TDDFT

## CO release of $Cr(CO)_6$ upon photoexcitation



Credit: J. Michl and V. Bonacic-Koutecký, "Electronic Aspects of Organic Photochemistry". John Wiley and Sons, Inc., New York, 1990.

#### $Cr(CO)_6$ - CO photodissociation

Experimentally, it is well-known that light-irradiation of carbonyl organometallic complexes leads to dissociation of a CO ligand and subsequent replacement by a solvent molecule.

$$Mo(CO)_6 + THF \xrightarrow{light} Mo(CO)_5 THF + CO$$

This mechanism is used in the lab to produce interesting initial reactants for more complex reactions.

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Study of  $Cr(CO)_6$  photodissociation using trajectory surface hopping to understand the electronic excited states involved.

## $Cr(CO)_6$ - CO photodissociation

#### Study of Cr(CO)<sub>6</sub> photodissociation using trajectory surface hopping

#### Computational details

- Gas phase system
- LR-TDDFT/B3LYP for 5 lowest excited state
- TSH dynamics
- $\bullet$  ~ 30 trajectories
- $\circ$   $\sim$  200 fs per trajectory

JCP, **134**, 164305 (2011)



#### $Cr(CO)_6$ - Comparison between methods

The root mean square deviation between experimental and theoretical excitation energies and oscillator stengths for two selected states.



#### CASPT2: importance of the active space.

#### $Cr(CO)_6$ - Potential energy surfaces

#### Validation of the potential energy surfaces



#### $Cr(CO)_6$ - Trajectory surface hopping

General trends of the nonadiabatic trajectories



Three steps: Hops within the large density of states - bands splitting - dynamics of the  $Cr(CO)_5$  fragment and relaxation towards  $S_0$ .

#### $Cr(CO)_6$ - Trajectory surface hopping

#### Different observables along the dynamics



## $Cr(CO)_6$ - Trajectory surface hopping

#### **Comparison with experimental findings**

TABLE V. Summary of simulated and experimental time constants (fs) for photodynamics of  $Cr(CO)_6$ . Experimental data from Ref. 16 for 345 nm pump wavelength. See Fig. 8-right for the definition of the time constants.

	Present work	Experimental
$t_d$	8	
te	11	
$\tau_1 = t_d + t_e$	19	$21 \pm 1$
$\tau_2 = t_d + 3t_e$	41	$43 \pm 3$
τις	148	$68 \pm 6$

 $t_d$  is the latency time before any dissociation takes place and  $t_e$  is the exponential time constant

 $\tau_1$ : Time when the dissociation starts.  $\tau_2$ : Time when Cr(CO)<sub>5</sub> population dominates (3 $t_e$  corresponds to 95% of the exponential decay distribution).  $\tau_{IC}$ : Excited state lifetime.

Nonadiabatic (Molecular) Dynamics