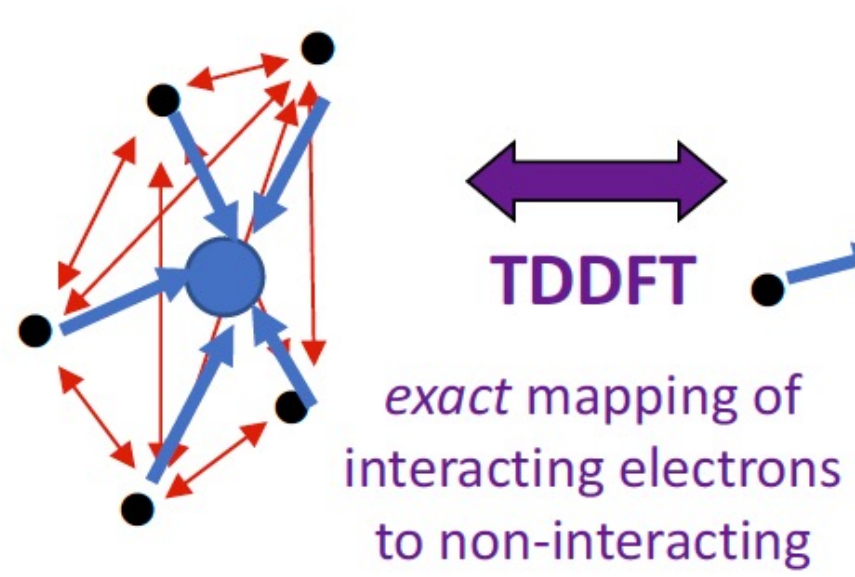


Theoretical Methods for Dynamics of Correlated Electrons, Ions and Photons

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Electron Dynamics via TDDFT

Time-dependent density functional theory (TDDFT) is the most widely used method to get electronic excitations and dynamics in molecules and solids.



Our work:

- Build improved approximations for the exchange-correlation potential (v_{XC}), especially for cases where the usual approximations fail.

$$i \partial_t \phi_i(\mathbf{r}, t) = (-\nabla^2/2 + v_s(\mathbf{r}, t)) \phi_i(\mathbf{r}, t)$$

$$v_s[n; \Phi_0](\mathbf{r}, t) = v_{ext}(\mathbf{r}, t) + \int d^3r' \frac{n(\mathbf{r}', t)}{|\mathbf{r} - \mathbf{r}'|} + v_{XC}[n; \Psi_0, \Phi_0](\mathbf{r}, t)$$

In particular:

Dependence of v_{XC} on the density $n(\mathbf{r}, t)$ at earlier times and the initial states.

Memory is missing from the usual approximations

Inaccurate dynamics (e.g.: charge-transfer, resonant driving, scattering...)

Our density-matrix coupled approach has memory, and satisfies exact conditions important for time-dynamics:

$$i \partial_t \rho_j(\mathbf{r}, t) = (-\nabla^2/2 + v_{ext}(\mathbf{r}, t) + v_H(\mathbf{r}, t) + v_{XC}(\mathbf{r}, t)) \rho_j(\mathbf{r}, t)$$

$$\nabla \cdot (n \nabla v_{XC}) = \nabla \cdot \left[\frac{1}{4} (\nabla' - \nabla) (\nabla^2 - \nabla'^2) (\rho_1(\mathbf{r}', \mathbf{r}, t) - (\rho_{1,s}(\mathbf{r}', \mathbf{r}, t)))|_{\mathbf{r}'=\mathbf{r}} + n(\mathbf{r}, t) \int n_{XC}(\mathbf{r}', \mathbf{r}, t) \nabla w(|\mathbf{r}' - \mathbf{r}|) \right]$$

$$i \frac{\partial \rho_1}{\partial t} = [-\nabla^2/2 + v_{ext}(\mathbf{r}, t)] \rho_1(\mathbf{r}, t) + \int d^3\bar{r} (w(\mathbf{r}, \bar{r}) - w(\mathbf{r}', \bar{r})) \rho_2[\rho_1, \rho_1](\mathbf{r}, \bar{r}; \mathbf{r}', \bar{r})$$

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L. Lacombe, N. T. Maitra, J. Chem. Theory and Comput. **15**, 1672 (2019).
L. Lacombe and N. T. Maitra, Faraday Discussions **224**, 382 (2020).

Coupled Electron-Ion Dynamics via EF

Dynamics of molecules beyond the Born-Oppenheimer approximation is increasingly relevant in many areas, e.g. photo-induced processes, solar-cell design, biomolecular mechanism modeling.

The Exact Factorization (EF) approach allows to understand how quantum subsystems' dynamics influence each other, and to develop mixed quantum-classical methods.

$$\hat{H} = \hat{H}_{BO} + V_{ext}^e(\mathbf{r}, t) + T_n(\mathbf{R}) + V_{ext}^n(\mathbf{R}, t)$$

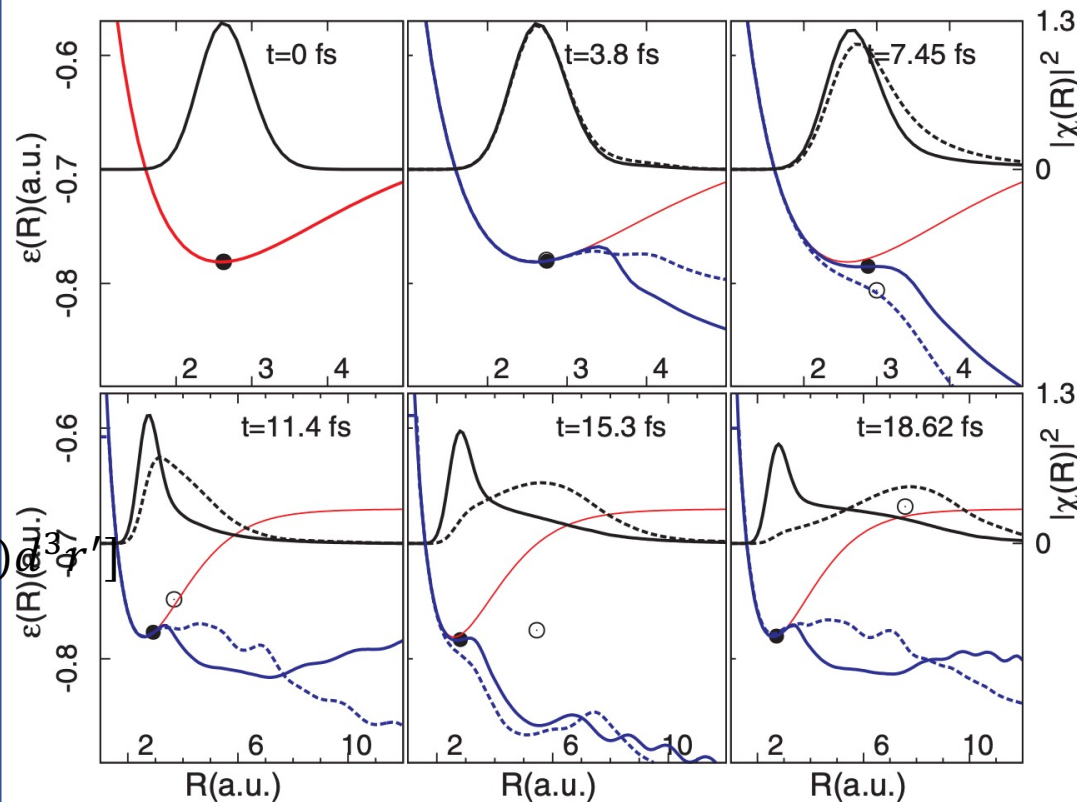
$$\hat{H}_{BO} = T_e(\mathbf{r}) + W_{ee}(\mathbf{r}) + V_{en}(\mathbf{r}, \mathbf{R}) + W_{ee}(\mathbf{R})$$

$$\hat{H}\Psi(\mathbf{r}, \mathbf{R}, t) = i \partial_t \Psi(\mathbf{r}, \mathbf{R}, t)$$

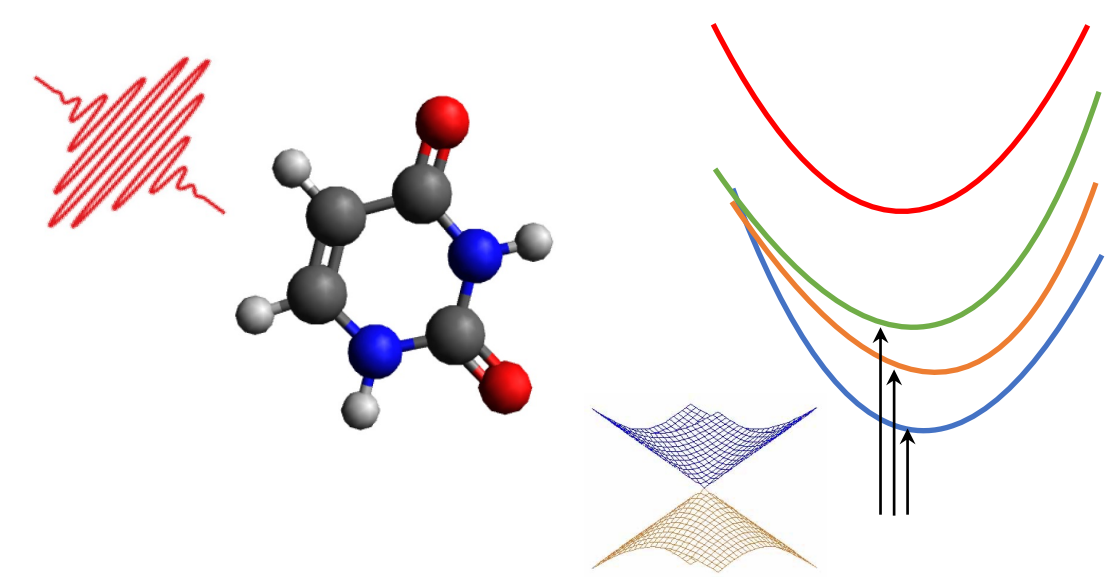
$$\Psi(\mathbf{r}, \mathbf{R}, t) = \Phi_{\mathbf{R}}(\mathbf{r}, t) \chi(\mathbf{R}, t) \text{ where } \int d\mathbf{r} |\Phi_{\mathbf{R}}(\mathbf{r}, t)|^2 = 1$$

The exact equations for χ and $\Phi_{\mathbf{R}}$ contain potentials that exactly capture the electron-nuclear correlation. They give rise to rigorous mixed quantum-classical methods, including surface-hopping approaches with first-principles decoherence corrections.

Exact potential driving electron in laser driven molecular dissociation



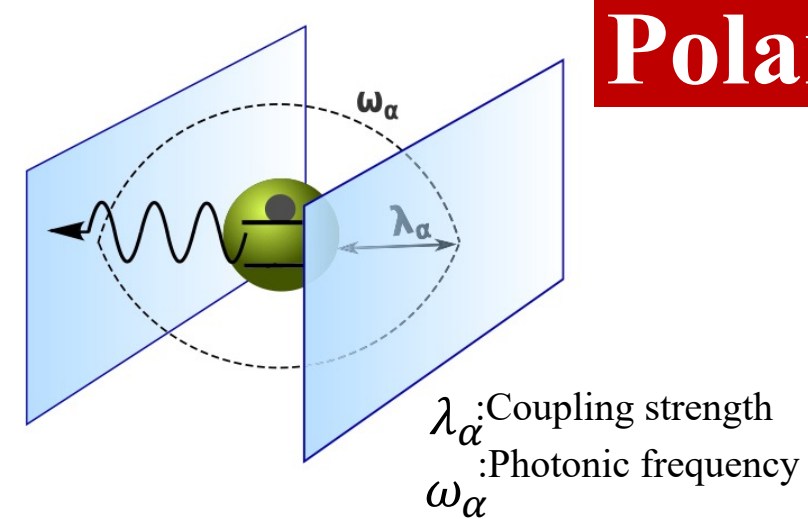
Coupled-trajectory approach based on the exact electron-nuclear correlation from the EF to describe quantum (de-)coherence in large molecules



Dynamics of excited uracil cation: quantum (de-)coherence can be treated correctly at the same computational cost as the original SH dynamics.

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G. H. Gossel, F. Agostini, N. T. Maitra, J. Chem. Theory Comput. **14**, 4513 (2018).
P. Vindel Zandbergen, L. M. Ibele, J.-K. Ha, S.-K. Min, B. F. E. Curchod, J. Chem. Theory Comput. (2021)

Polaritonic Chemistry



Molecule in a nanoscale cavity greatly enhances the light-matter coupling strength

Hybrid light-matter states (polariton)
Quantized nature of light becomes important

Our work:

Exact factorization and trajectory methods to gain insight and model the new phenomena

Manipulate light and matter properties (recent experimental advances in cavity QED)

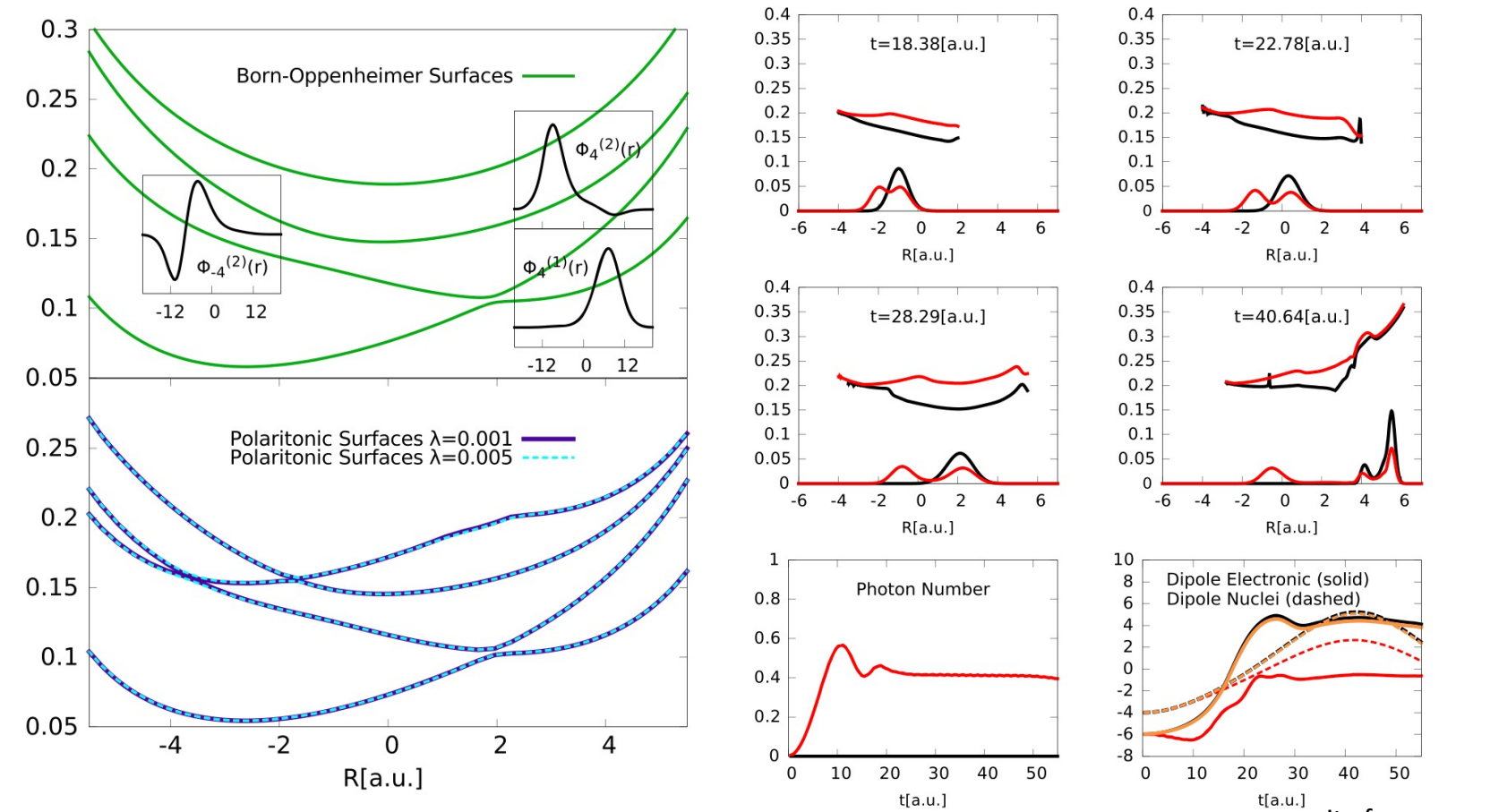
$$\Psi(\mathbf{r}, \mathbf{R}, q, t) = \chi(q, t) \Phi_q(\mathbf{r}, \mathbf{R}, t) = \chi(\mathbf{R}, t) \Phi_{\mathbf{R}}(\mathbf{r}, q, t) = \chi(\mathbf{r}, t) \Phi_{\mathbf{r}}(q, \mathbf{R}, t)$$

- TDSE for the photons
- corrections to the quadratic form due to matter-photon coupling

- TDSE for the nuclei
- Exact potential driving the nuclei coupled to electrons and photons

EF gives cavity-modified time-dependent potential energy surfaces for proton-coupled electron transfer that directly indicate suppression:

Ehrenfest trajectories for photons' displacement field coordinates work well.



L. Lacombe, N. M. Hoffmann, N. T. Maitra, Phys. Rev. Lett. **123**, 083201 (2019).
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Electronic Embedding via EF (EVEF)

Systems containing strong correlation require us to go beyond usual density-functional methods. We are developing a new, practical, first-principles quantum electronic embedding based on exact factorization (EF) in Fock Space:

For $H\Psi = E\Psi$, factorize

$$\Psi(n_1, n_2, \dots, n_k, n_{k+1}, \dots, n_M) = \chi(\tilde{n}) \Phi_{\tilde{n}}(\tilde{m})$$

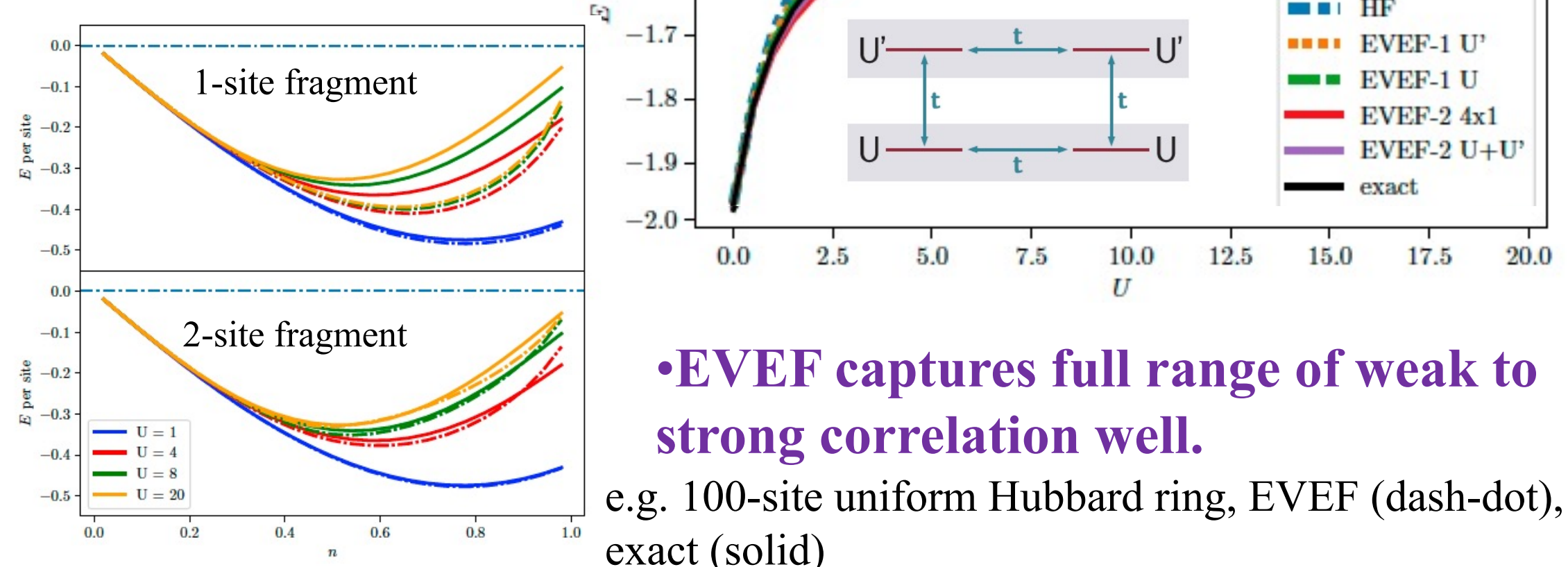
Site-occupations for chosen fragment, e.g. those with strongly-correlated orbitals

$$\sum_m \Phi_m^*(\tilde{m}) \Phi_m(\tilde{m}) = 1$$

Embedding Hamiltonian: $h_{\tilde{n}, \tilde{n}'} \equiv \sum_{m', m} \Phi_m^*(\tilde{m}) H_{n, m; n', m'} \Phi_{m'}(\tilde{m}')$

Use mean-field method $\rightarrow \Psi \rightarrow$ extract $\Phi_{\tilde{n}}$ and embedding $h \rightarrow$ solve $h\chi = E\chi$ for fragment χ with high-level method \rightarrow accurate energy E

e.g. Hubbard tetramer: two sites with fixed interaction $U' = 0.1$ and two sites with variable U



EVEF captures full range of weak to strong correlation well.
e.g. 100-site uniform Hubbard ring, EVEF (dash-dot), exact (solid)

L. Lacombe and N. T. Maitra, Phys. Rev. Lett. **124**, 206401 (2020)



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